



Wir schaffen Wissen – heute für morgen

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Use of SiPMs in ZnS:⁶LiF scintillation neutron detectors



- widely used in the neutron scattering experimental technique
- currently the only real alternative to the 3He gas detectors
- 1D or 2D spatial resolution
- number of detection channels: $10^3 10^4$
- covered area: $1 10 \text{ m}^2$

<u>Thermal neutrons</u>: E = 1 - 100 meV ($\lambda = 1 - 6 \text{ Å}$)

Detection:

1. conversion into charged particles (triton + alpha) through the nuclear absorption reaction on the ⁶Li isotope (neutron converter):

 $^{6}\text{Li} + {}^{1}\text{n} \rightarrow {}^{3}\text{H} + {}^{4}\text{He} + 4.79 \text{ MeV} \qquad \sigma = 940 \cdot \lambda / 1.8 \text{ [barn]}$

2. detection of the products in the ZnS(Ag) scintillator

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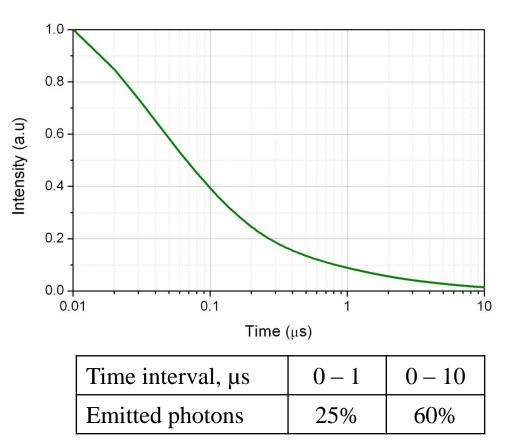
Neutron detection screens from Applied Scintillation Technologies http://www.appscintech.com

	ND4:1	ND2:1	
Mass ratio ZnS:6LiF	4:1	2:1	
Density, g/cm ³	2.2	2.2	
⁶ Li atoms, 10 ²² cm ⁻³	1.0	1.5	
Thickness, mm	0.45, 0.25		
Emission max., nm	450		
Photons per neutron	160000		
Transparency	opaque		

- bright (+)
- non-transparent (-)
- usable thickness ≤ 0.5 mm (-)
- scintillation process slow (-)

ND scintillator luminescence in response to neutrons (from E.S.Kuzmin et.al., Journal of Neutron Research 10 (2002) 31)

Ampl	191	230	88	50	25	6	1.2
τ, μs	0.022	0.074	0.208	0.88	4.3	18.1	87.7



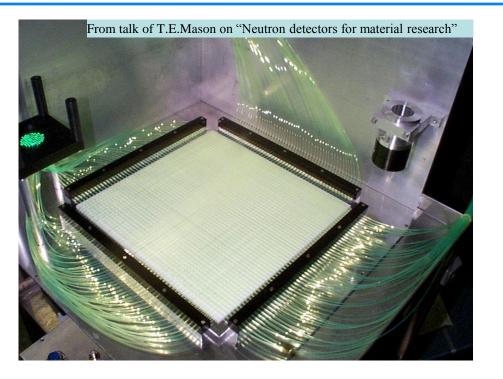




- V-shape (or Venetian blind) arrangement of
 the scintillator stripes to increase the
 neutron path through the scintillator and
 thus the probability of its absorption
- light collection through clear fibers
- channel coding by sharing the light between several PMTs (restricts the number of PMTs)

No way for using SiPMs, large-area photosensor is required!





- light collection by 1D(2D) arrays of WLS fibers
- light detection by MaPMTs

Direct (without design change) replacement of MaPMTs by SiPMs is possible, but the dark count rate of the SiPM has to be substantially reduced \rightarrow requires deep cooling.

The necessity to cool SiPMs will add to the complexity and to the price of the detector, which might otweigh all the advantages expected from their usage.

Our goal: a WLS fiber based detector readout by SiPMs operated at room temperature.



Detection efficiency at 1Å, %	40 - 50	
• neutron absorption probability at 1Å, %	50 - 60	
• trigger efficiency, %	70 - 90	Problematic to fulfill simultaneously
Background count rate, Hz	< 10 ⁻³	with a SiPM operated at RT (dark count rate ~ 1 MHz)
Gamma-sensitivity (probability to detect 1.2 MeV γ from ⁶⁰ Co source)	< 10 ⁻⁶	
Dead time, µs	5 - 20	
Multi-count ratio, % (mean number of triggers per event – 1)	< 1	

To be able to use SiPMs at RT we need:

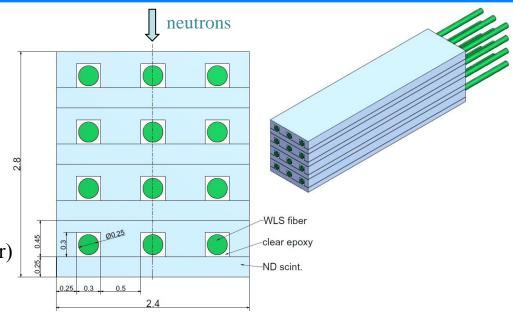
- 1. a much better light collection from the scintillator (at least a factor of 10)
- 2. a signal processing system capable to reliably identify the neutron signals against the high background of the SiPM dark counts

Single-channel detection unit (pixel in a multichannel detector)

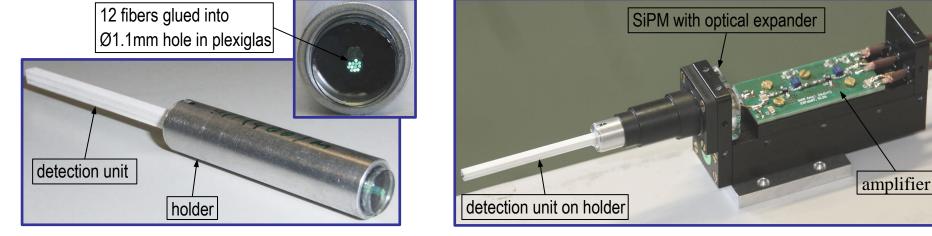
2.4 x 2.8 x 50mm

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- effective absorption volume = 0.83
- 12 fibers Y11(400)M, Ø = 0.25mm
- neutron absorption probability: ~ 80% at 1Å
- readout with 1x1mm² SiPM via "optical expander" (Ø1.2x5mm clear fiber)

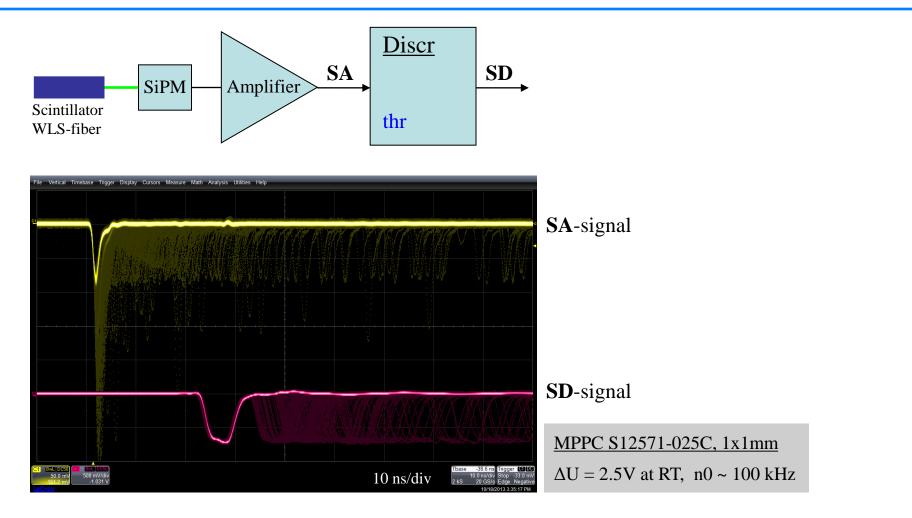


Single-channel prototype detector for current measurements

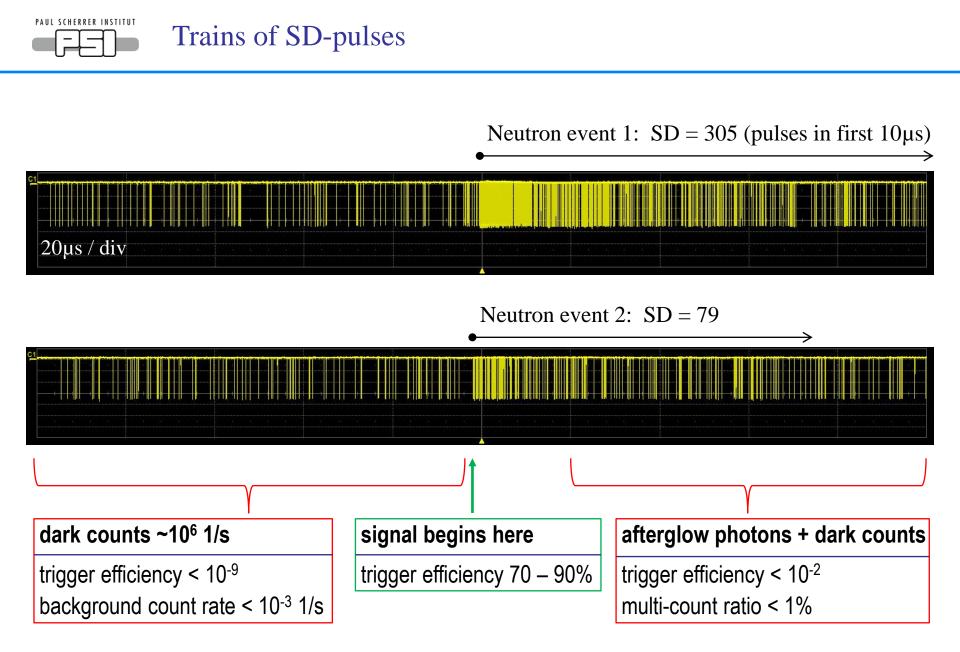




Signal processing (step 1): analog \rightarrow digital

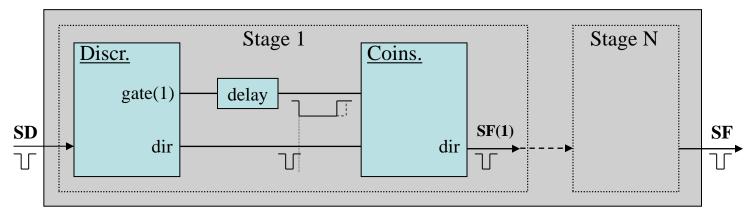


Independent of how many SiPM cells fire at the same time, only one SD-pulse is generated per one primary electron. Cross-talk is completely removed, further signal processing is independent of the used SiPM type.



Removal of dark counts / Signal extraction

Multistage single-pulse elimination filter (high-pass pulse density filter)



SD - input pulse sequencegate(i) - width of the self-coincidence gate for stage(i)SF - subsequence of SD passing the filterN - number of stages

To pass the stage(i) a pulse of the input sequence should have a preceding pulse within the time interval gate(i). Pulses not satisfying this condition are removed.

Dark counts suppression at gate(1) = gate(2) = ... = gate(N):

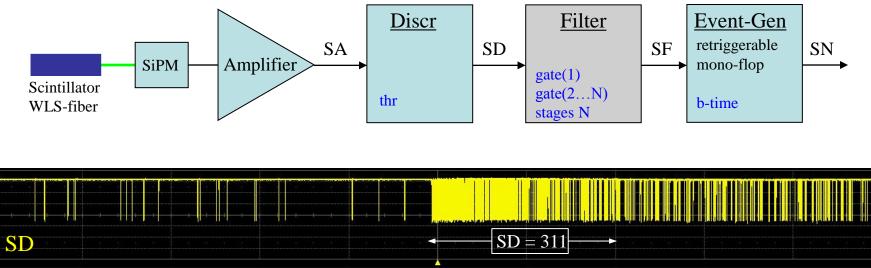
bgd = **n0** (1 – α) α^{N} , $\alpha = \Delta n0$, $\Delta = gate(1) + SD$ -pulse width, n0 - SiPM dark count rate

Practical choice of gates: $gate(1) \ll gate(2) = ... = gate(N)$

- -- dark count suppression almost as above, mainly determined by gate(1)
- -- better transmission of neutron signals



Signal-processing: full chain

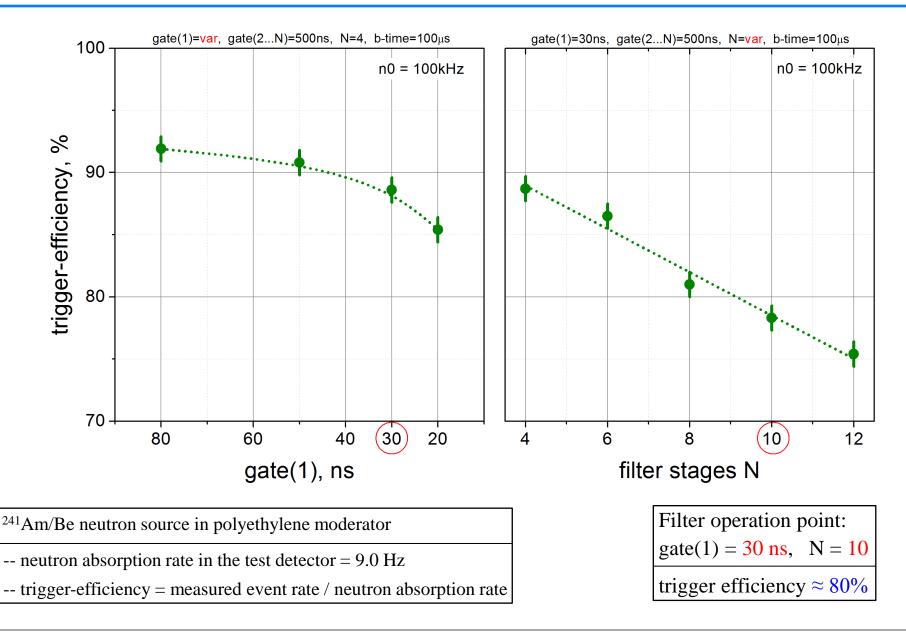




A retriggerable mono-flop as event generator \implies "blocking time" is automatically adjusted to the strength of SF-signal.

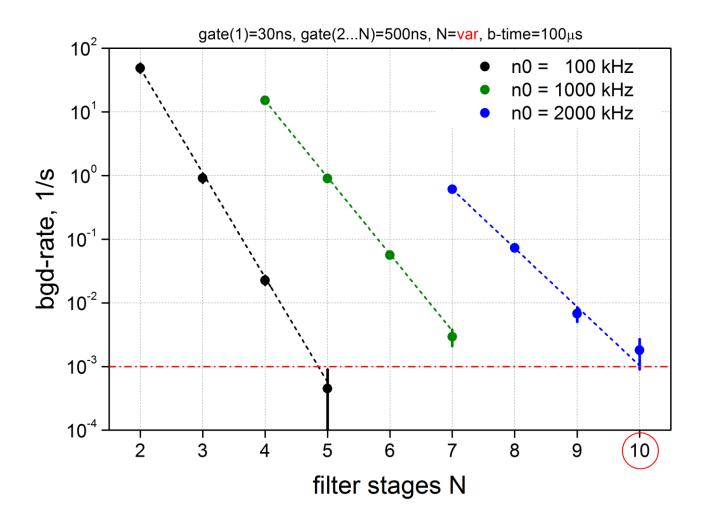
Further measurements	parameter	thr	gate(1), ns	gate(2N), ns	stages N	b-time, μs
with the following settings \implies	value range	$0.5 A_{1e}$	20 - 80	500	4 - 12	5 - 100







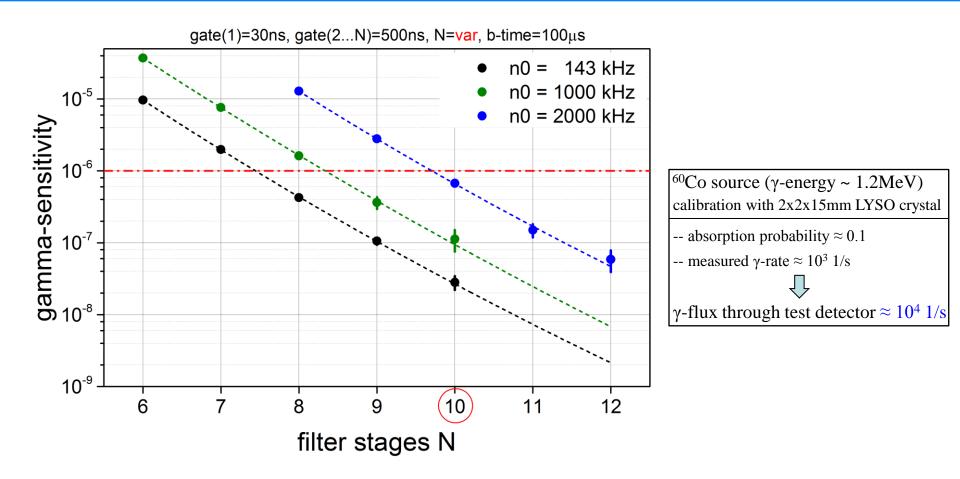
Background count rate



At trigger-efficiency of ~ 80% the background count rate $\leq 10^{-3}$ Hz is guarantied up to the SiPM dark count rate of ~ 2 MHz.

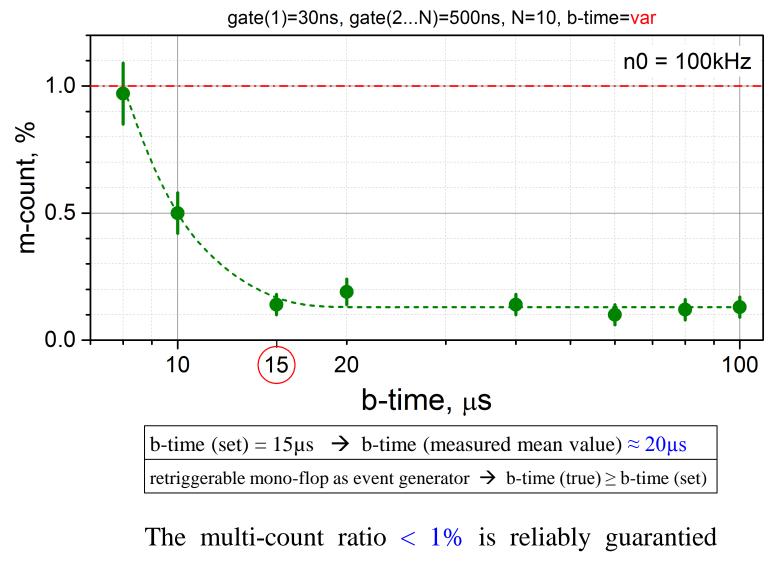


Gamma-sensitivity



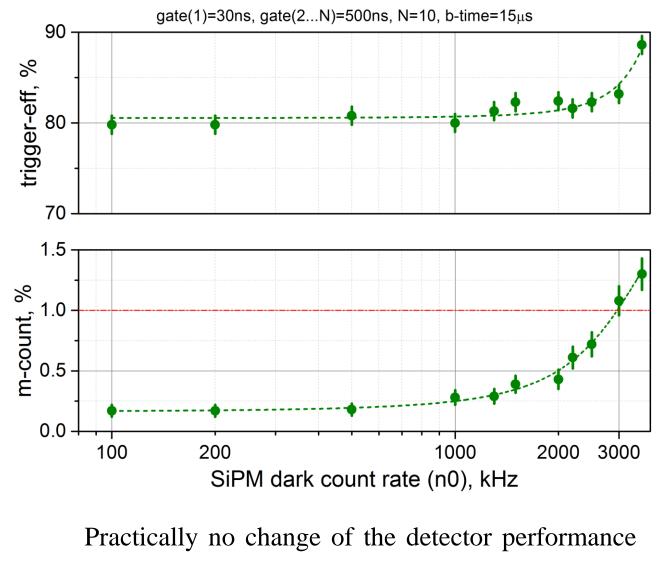
At trigger-efficiency of ~ 80% the gamma-sensitivity $\leq 10^{-6}$ is guarantied up to the SiPM dark count rate of ~ 2 MHz.





with the true blocking time of $\sim 20 \mu s$.

Multi-count ratio, efficiency vs. SiPM dark count rate



up to the SiPM dark count rate of ~ 2 MHz.

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We proved the feasibility of using Silicon Photomultipliers operated at room temperature for individual readout of detection channels in multichannel neutron detectors utilizing ZnS:⁶LiF scintillation screens.

The problem of the high dark count rate of the SiPM is solved by substantially improving the light collection from the scintillator and by developing an efficient algorithm based on the time-domain filtering for separating the signals from the dark counts.

With a single-channel prototype detector we achieved very good performance characteristics up to 2 MHz SiPM dark count rate.



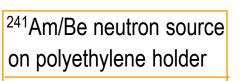
Summary 2

We are convinced, that the use of SiPMs in ZnS:⁶LiF scintillation detectors for application in the neutron scattering experimental technique will open on the one hand new possibilities in designing this kind of detection systems, and on the other hand a new field for the large-scale application of SiPMs.



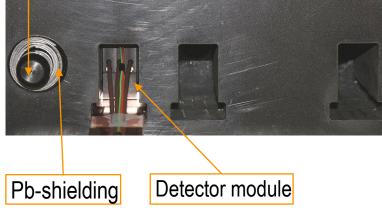


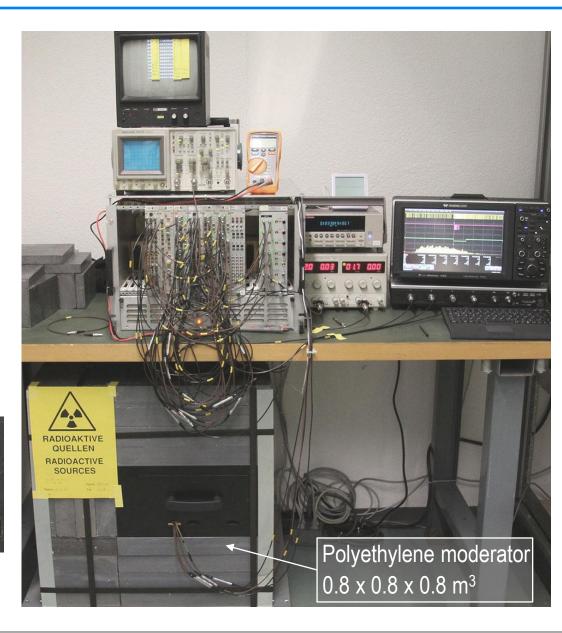
Measurement setup



 $2 \cdot 10^4$ fast neutrons / s

Moderator inner volume with caverns

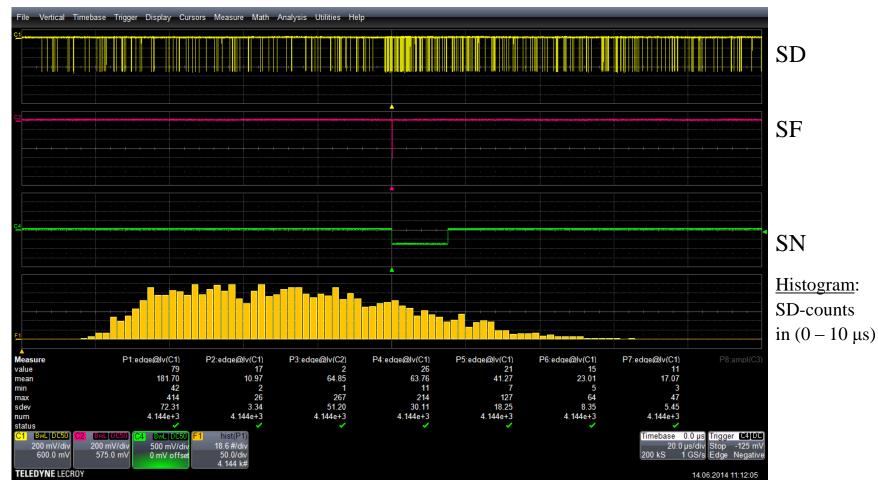






DSO event detection panel

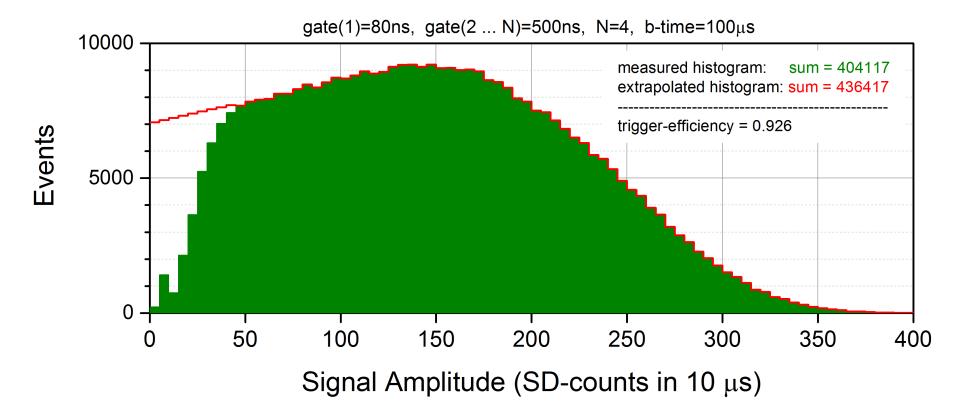
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Parameter	P1	P2	Р3	P4	P5	P6	P7
Window, µs	(-90, -80)	(0, 10)	(0, 10)	(10, 20)	(20, 30)	(50, 60)	(90, 100)
Description	SD-counts	SD	SF	SD	SD	SD	SD



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-- trigger efficiency by extrapolating the SD-amplitude spectrum to zero amplitude: 0.926 -- measured event rate: 8.3 Hz

neutron absorption rate in the scintillator: 9.0 Hz



 $E \text{ [meV]} = 81.82 / \lambda^2 \text{ [Å]}$ λ [Å] = 9.045 / E^{0.5} [meV] $v [m/s] = 3956 / \lambda [Å]$

<i>E</i> , meV	λ, Å	v, m/s	$\Delta t(1 \text{cm}), \mu s$
81.8	1.0	3956	2.5
25.2	1.8	2197	4.5
2.3	6.0	659	15.7

 $\Delta t(1 \text{ cm})$ – travel time in 1 cm

Detection

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 $^{3}\text{He} + ^{1}\text{n} \rightarrow ^{3}\text{H} + ^{1}\text{p} + 0.76 \text{ MeV}, \quad \sigma = 5333 \cdot \lambda / 1.8 \text{ [barn]}$ $^{6}\text{Li} + {}^{1}\text{n} \rightarrow {}^{3}\text{H} + {}^{4}\text{He} + 4.79 \text{ MeV}, \quad \sigma = 940 \cdot \lambda / 1.8 \text{ [barn]}$

Interaction probability

 $\varepsilon = 1 - \exp(-N \cdot \sigma \cdot d)$

N [cm⁻³] – density of absorbing atoms

 σ [barn] – absorption cross-section

d – detector thickness

Density of absorbing atoms: ³He: $2.7 \cdot 10^{19} \text{ cm}^{-3} \cdot \text{atm}^{-1}$ ND2:1 scint: $1.5 \cdot 10^{22}$ cm⁻³

Attenuation length at 1Å ³He (1 atm): **12 cm** : 0.13 cm ND2:1